

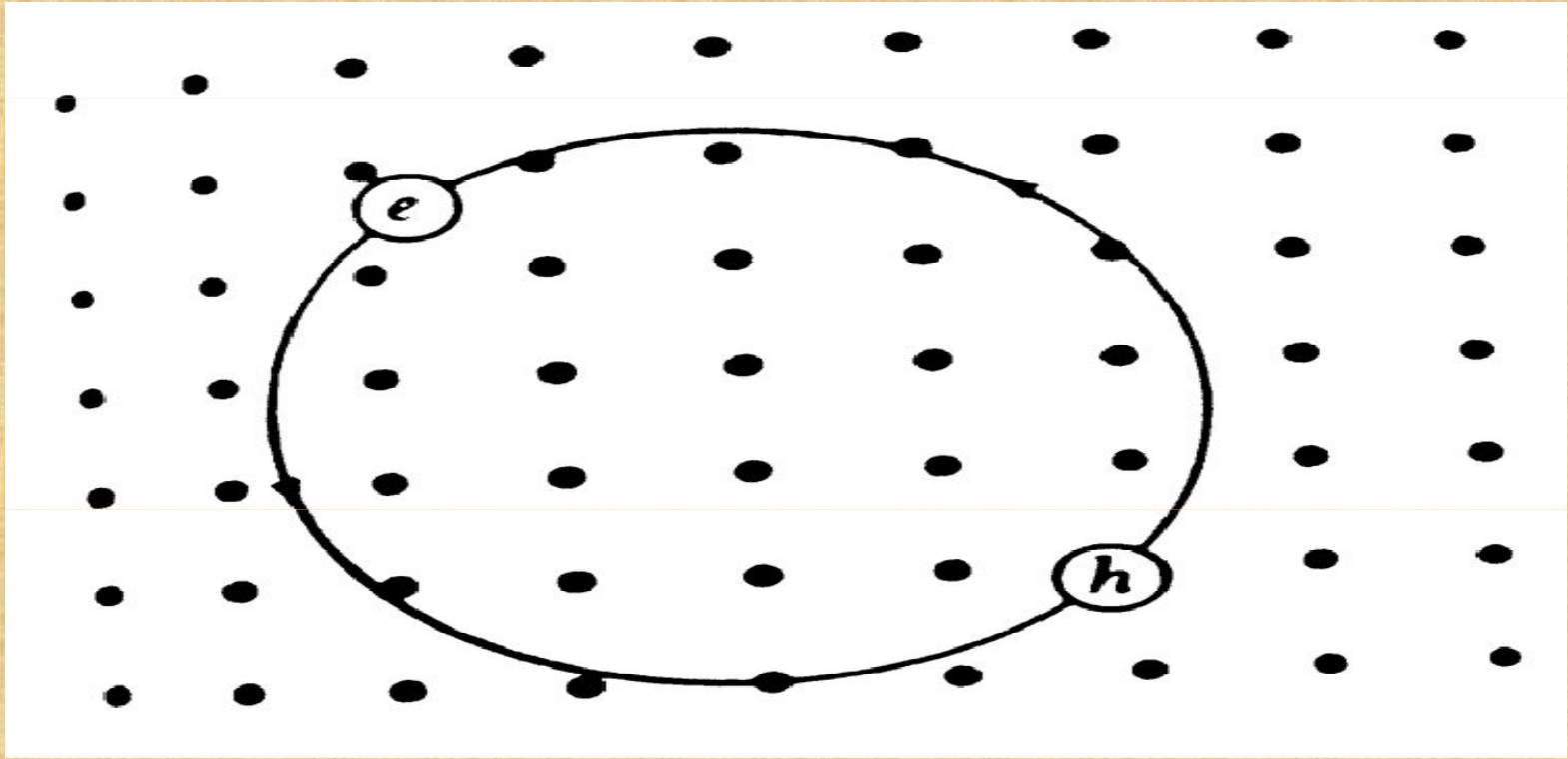
Excitons

excitons

- Frenkel Excitons
- Alkali Halides
- Molecular crystals
- Weakly bound (**Mott-wannier**) Excitons
- Exciton Condensation into Electron –Hole
- Drops(EHD)

EXCITONS

- An electron and a hole may be bound together by their attractive coulomb interaction , just as an electron is bound to a proton to form a neutral hydrogen atom . the bound electron-hole pair is called an exciton .
an exciton can move through the crystal and transport energy ; it does not transport charge because it is electrically neutral



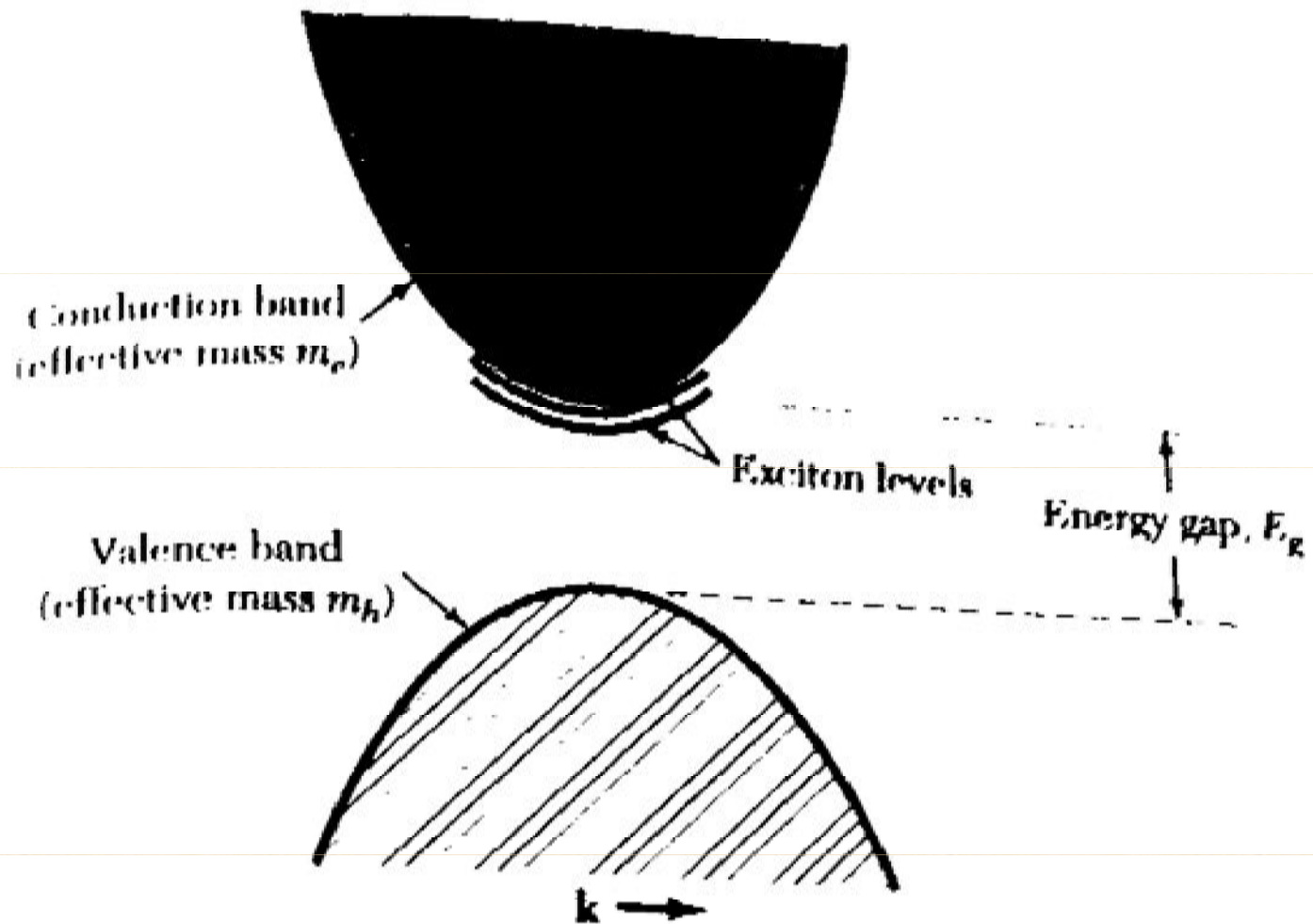
Characteristics of exciton

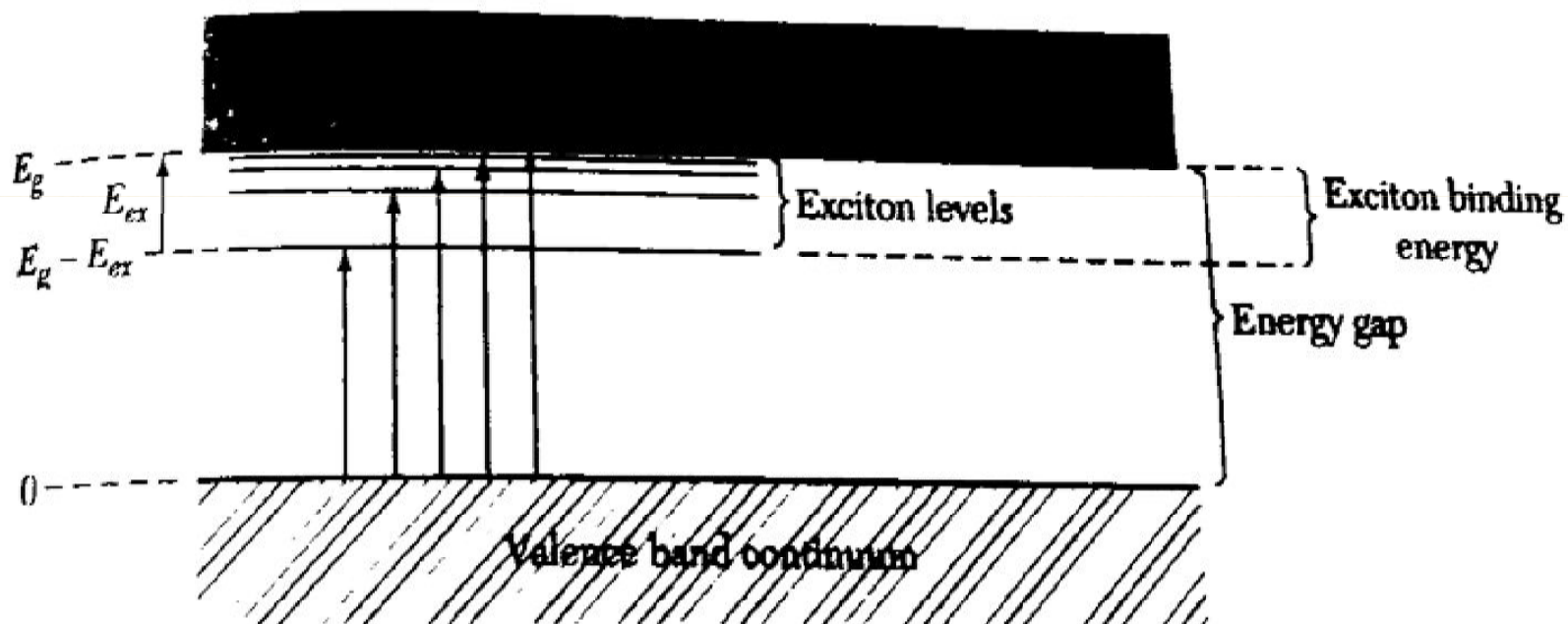
- It is formed in every insulating crystal
- It can form complexes;



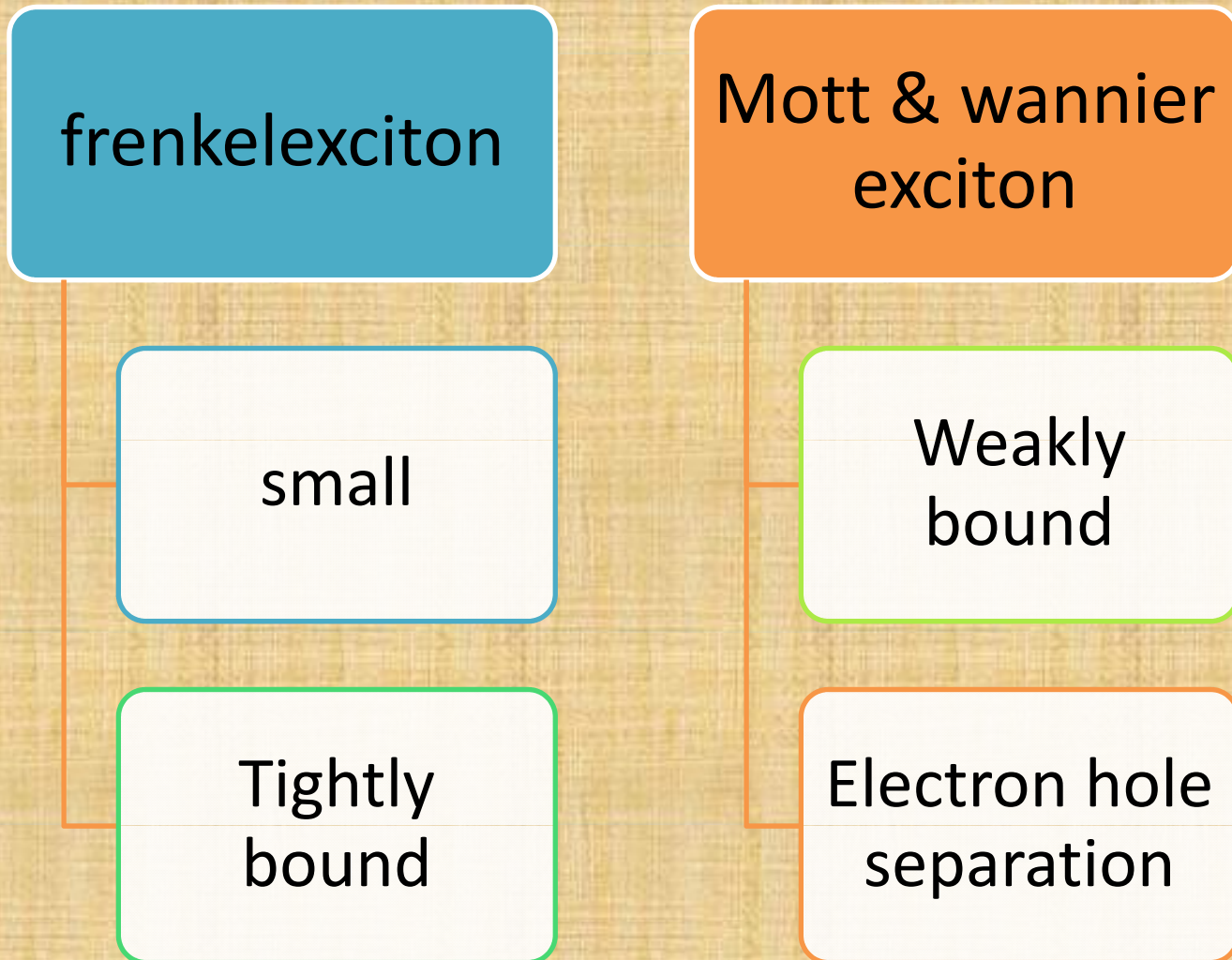
The binding energy of the exciton can be measured in three ways

- In optical transitions from the valence bond , by the difference between the energy required to create an exciton and the energy to create a free electron & free hole.
- In recombination luminescence, by comparison of the energy of the free electron hole recombination line with the energy of the exciton recombination line.
- By photo-ionization of excitons ,to form free carriers . this experiment requires high concentration of excitons .



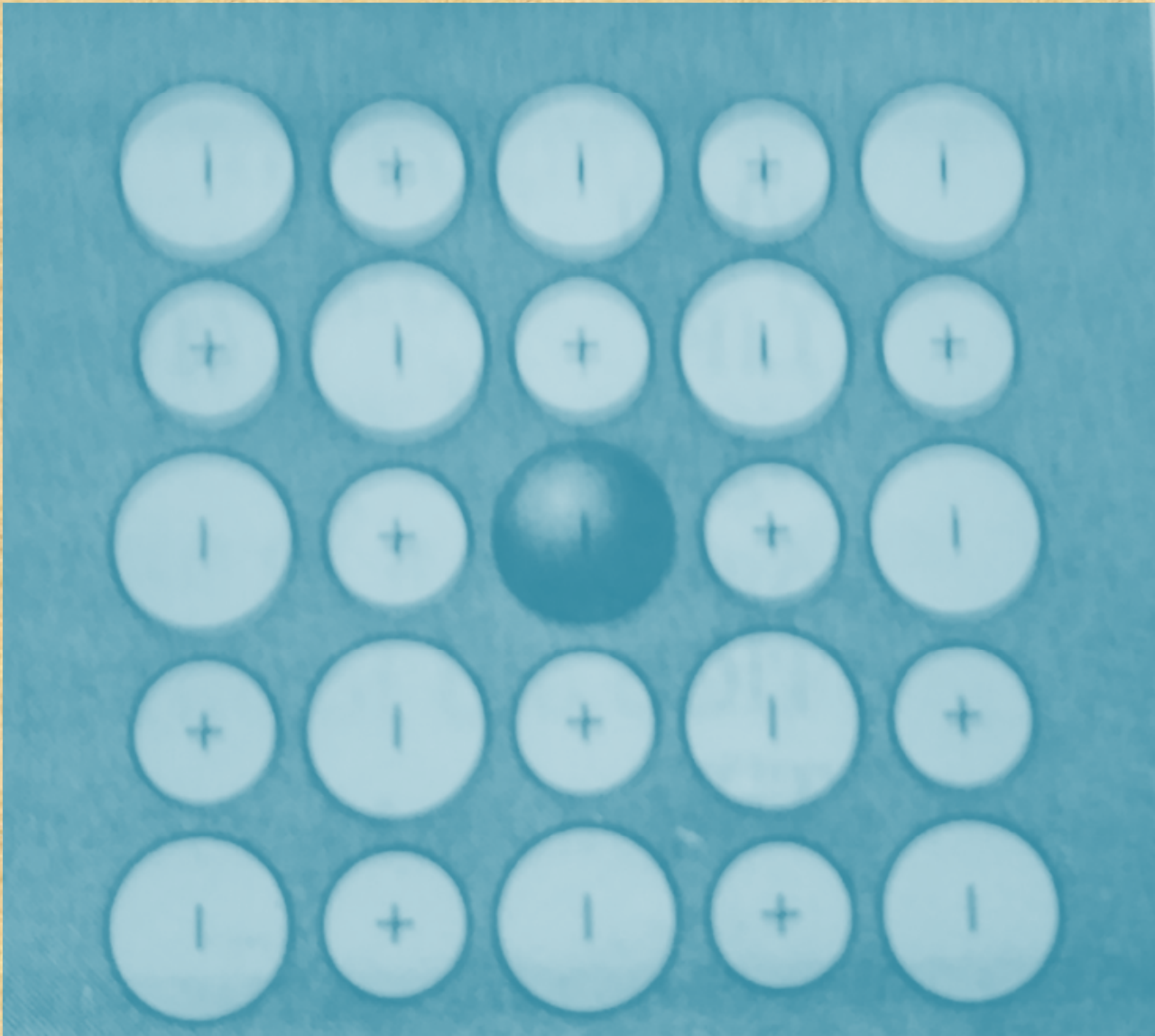


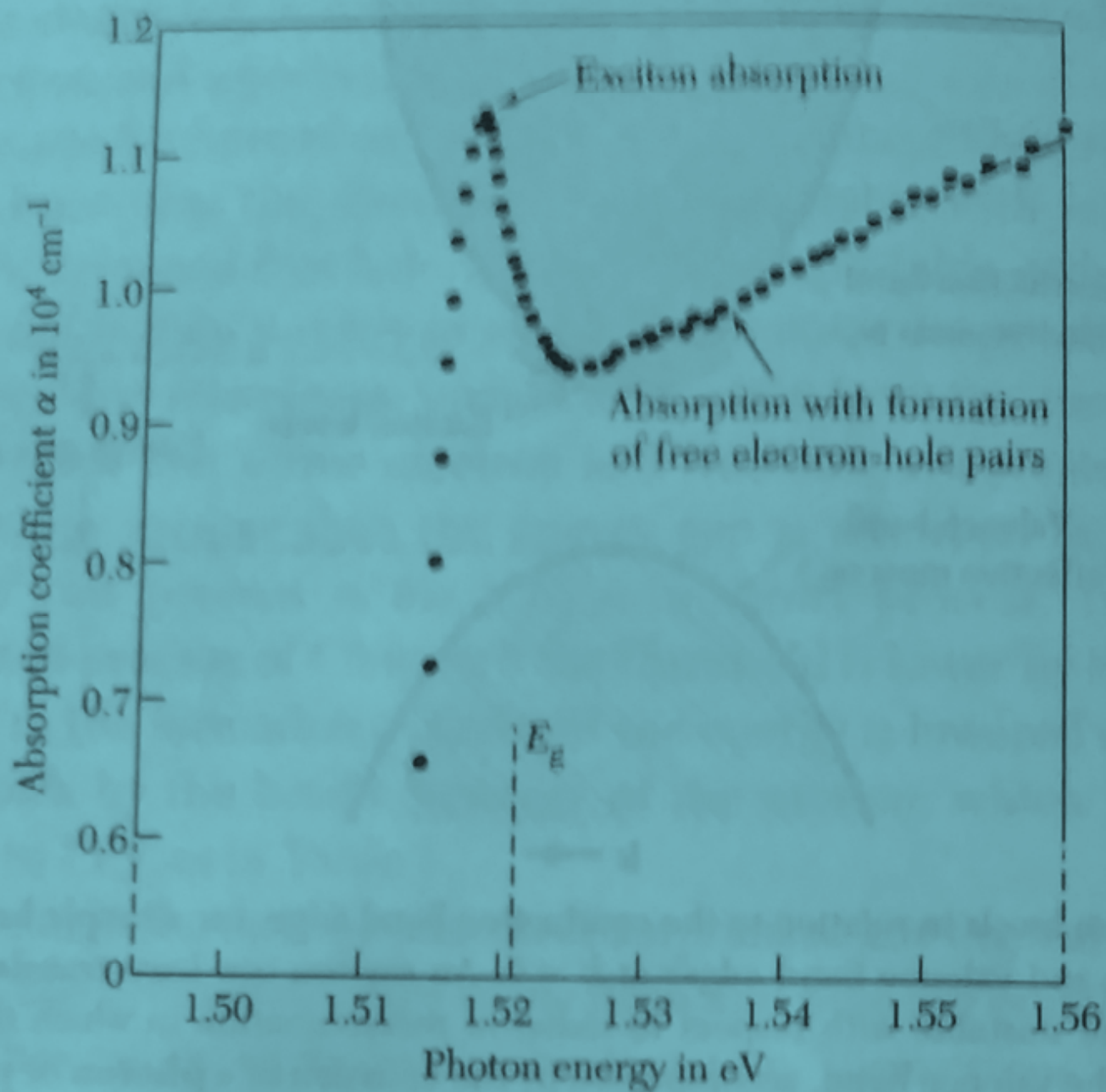
Excitons in two different approximations



Frenkel exciton

- An excited state of single atom
- Exciton can hop from one atom to another by coupling
- Exciton waves travels through the crystal much as the reversed spin of a magnon travels through the crystal.
- For example krypton.....,
exciton ground state energy= $11.7-10.17$
 $=1.5$ eV





$$\psi_g = u_1 u_2 \cdots u_{N-1} u_N ,$$

$$\varphi_j = u_1 u_2 \cdots u_{j-1} v_j u_{j+1} \cdots u_N .$$

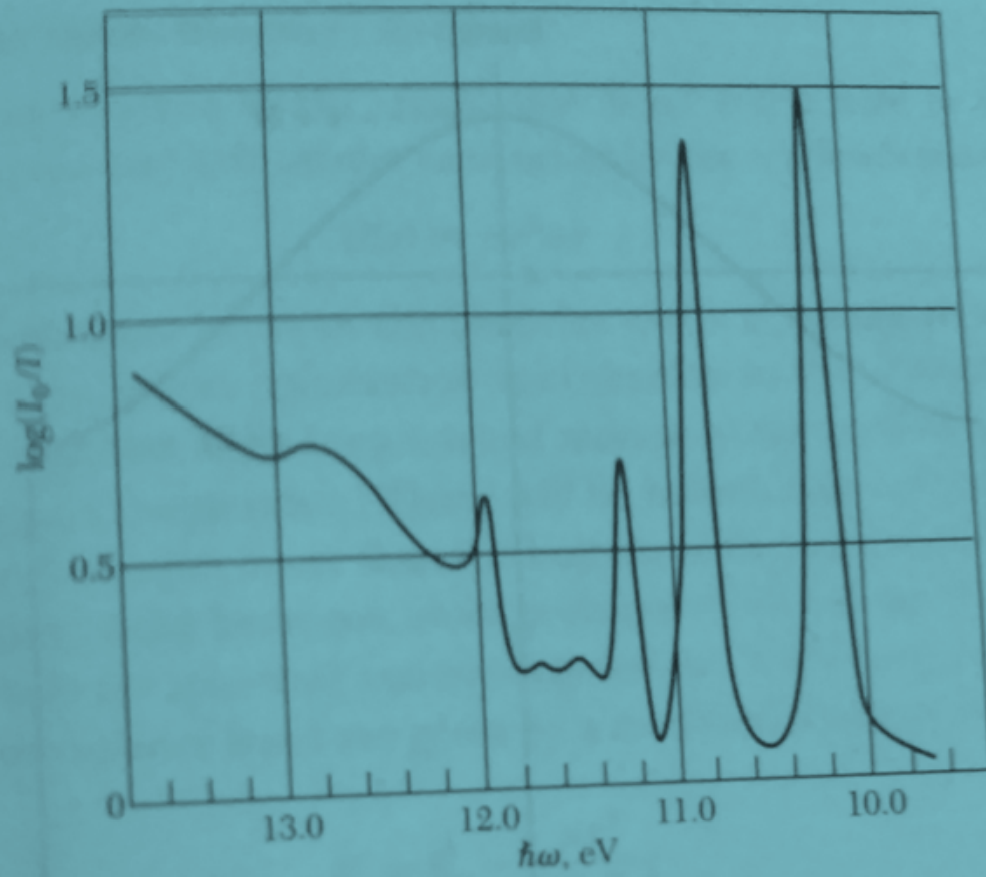


Figure 8 Absorption spectrum of solid krypton at 20 K. (After G. Baldini.)

$$\mathcal{H}\varphi_j = \epsilon\varphi_j + T(\varphi_{j-1} + \varphi_{j+1})$$

$$\psi_k = \sum_j \exp(ijka) \varphi_j .$$

To see this we let \mathcal{H} operate on ψ_k :

$$\mathcal{H}\psi_k = \sum_j e^{ijka} \mathcal{H}\varphi_j = \sum_j e^{ijka} [\epsilon\varphi_j + T(\varphi_{j-1} + \varphi_{j+1})] ,$$

We rearrange the right-hand side to obtain

$$\mathcal{H}\psi_k = \sum_j e^{ijka} [\epsilon + T(e^{ika} + e^{-ika})] \varphi_j = (\epsilon + 2T \cos ka) \psi_k ,$$

so that the energy eigenvalues of the problem are

$$E_k = \epsilon + 2T \cos ka ,$$

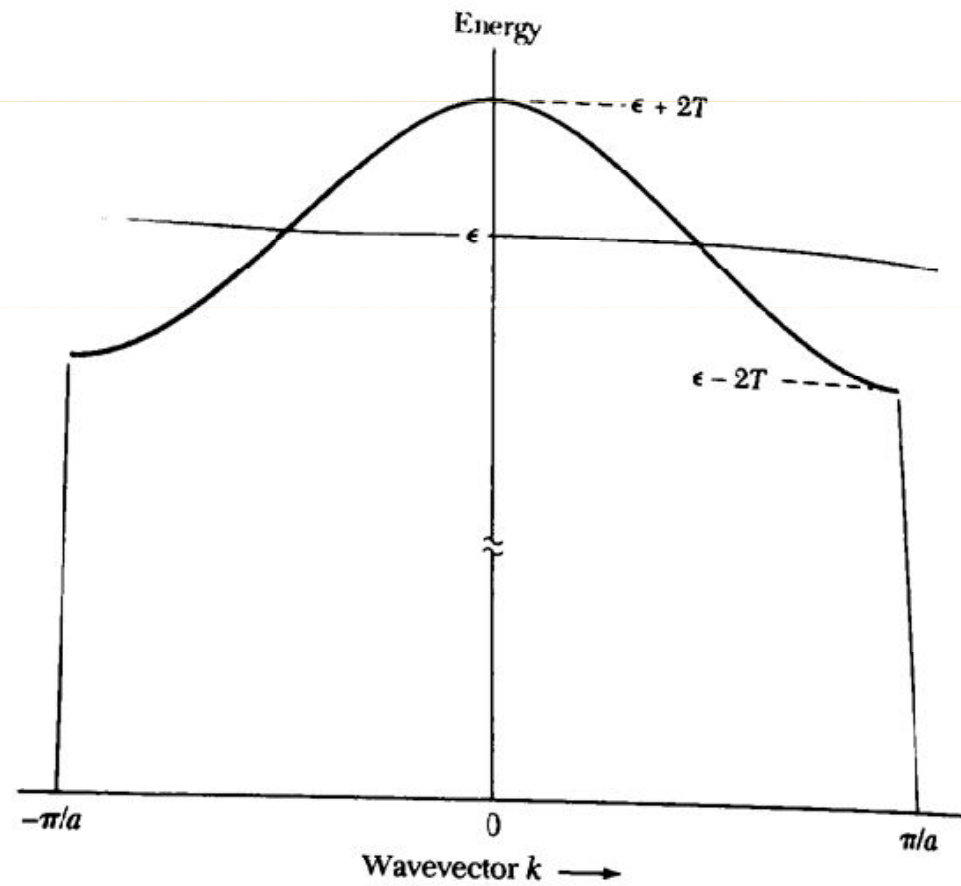


Figure 9 Energy versus wavevector for a Frenkel exciton, calculated with positive neighbor transfer interaction T .

$$k = 2\pi s/Na; \quad s = -\frac{1}{2}N, -\frac{1}{2}N + 1, \dots, \frac{1}{2}N - 1$$

Alkali halides

- Lowest energy excitons are localized on the negative halogen ions.
- Negative ions $<$ positive ions.
- transparent – in visible spectrum region.
- Energies do not lie in that region.
- Excitonic absorption structure in the vacuum ultra violet.
- Splitting is caused by spin orbit interaction.

Molecular crystals

- The covalent binding is strong within the molecule compared to vanderwaals binding between the molecules.
- Electronic excitation lines of an molecule appear in the crystalline solid with little shift in frequency.
- At low temperature the lines in the solid are sharp , but there will be more lines due to davydov splitting

Weakly bound excitons

- The electron & hole attract each other by the coulomb potential,

$$U(r) = -e^2/\epsilon r$$

$$E_n = E_g - \frac{\mu e^4}{2\hbar^2 \epsilon^2 n^2} .$$

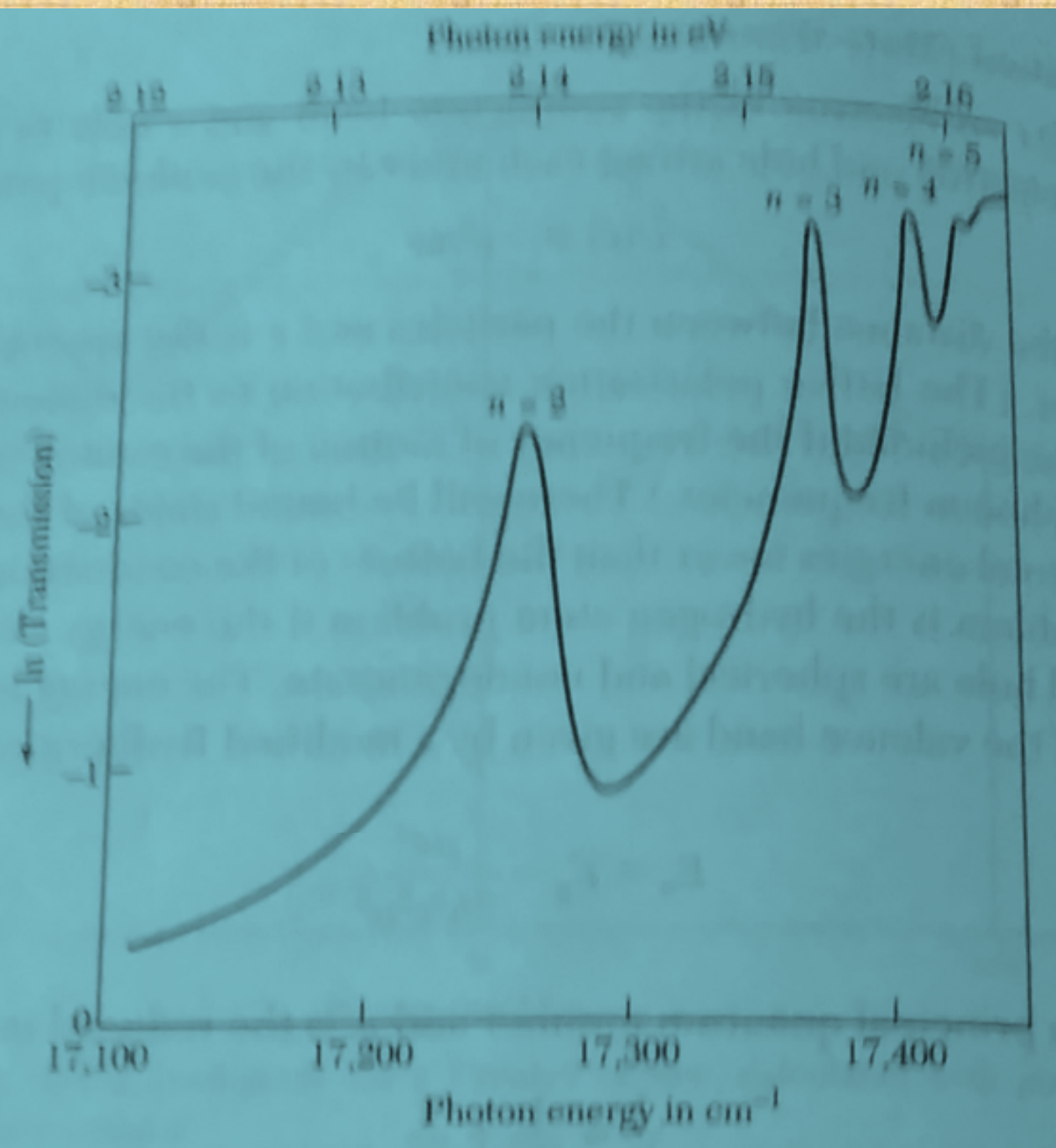
Here n is the principal quantum number and μ is the reduced mass:

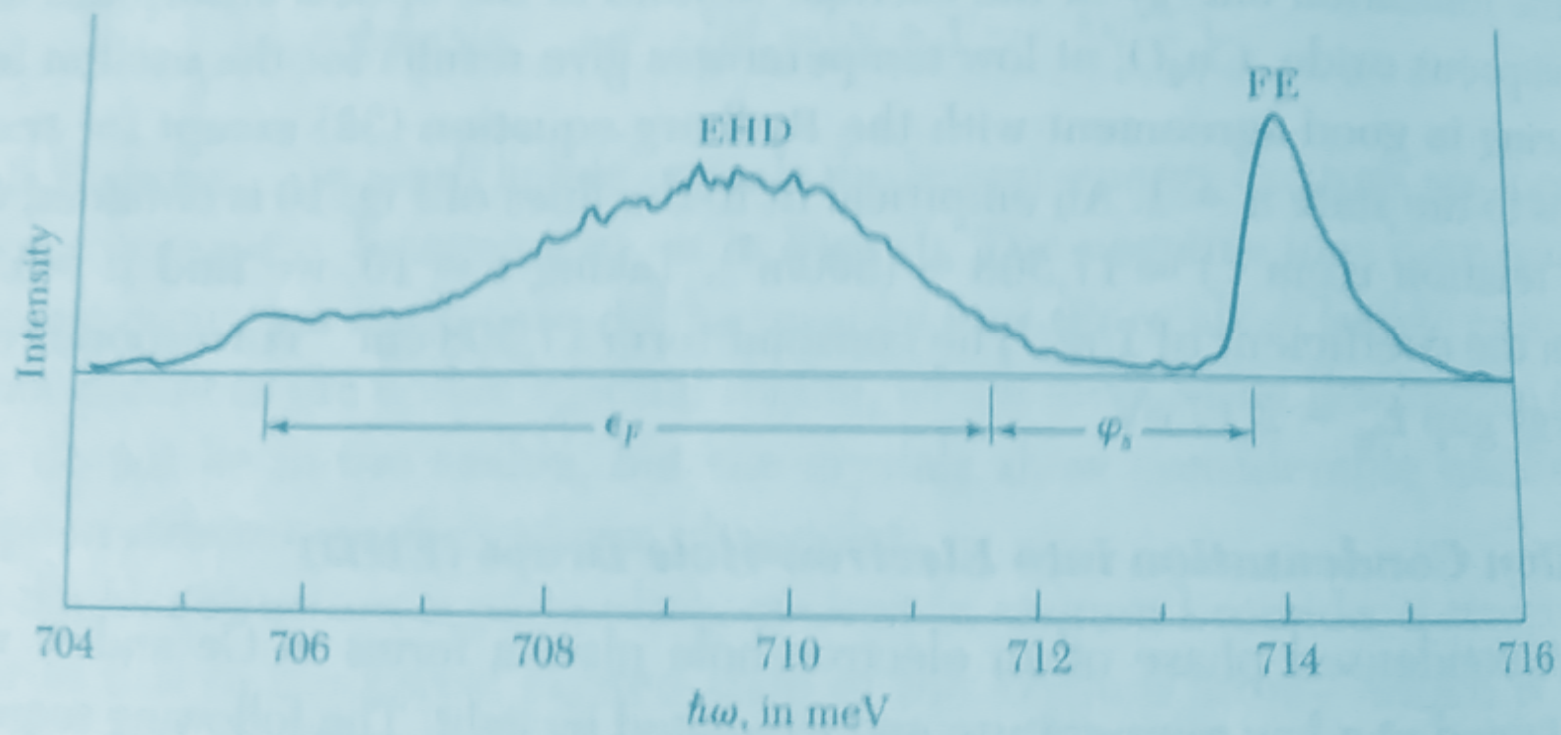
$$\frac{1}{\mu} = \frac{1}{m_e} + \frac{1}{m_h} ,$$

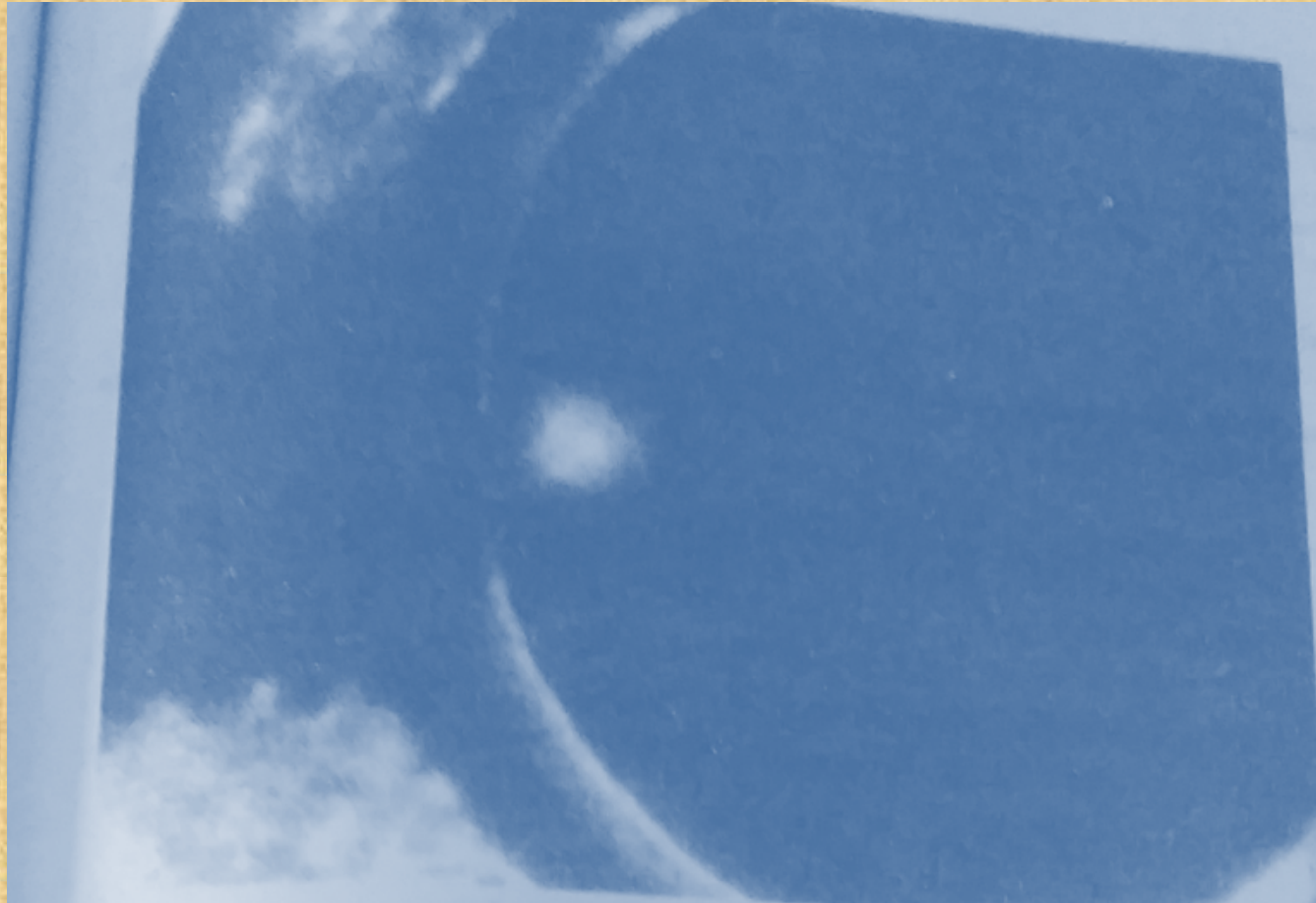
Taking $E=10$ & $\mu=0.7m$ from the coefficient of $1/n^2$. the constant term is $17,508 \text{ cm}^{-1}$,corresponding to an energy gap $E_g = 2.17 \text{ eV}$

Exciton condensation into electron hole drops (EHD)

- at low temperature & irradiated by light
Ge & Si .
- In Ge ,the absorption of photon of free energy produces a free electron and free hole with high efficiency.
- Lifetime -8 μs
- Drop lifetime -40 μs
- In strained Ge as long as 600 μs







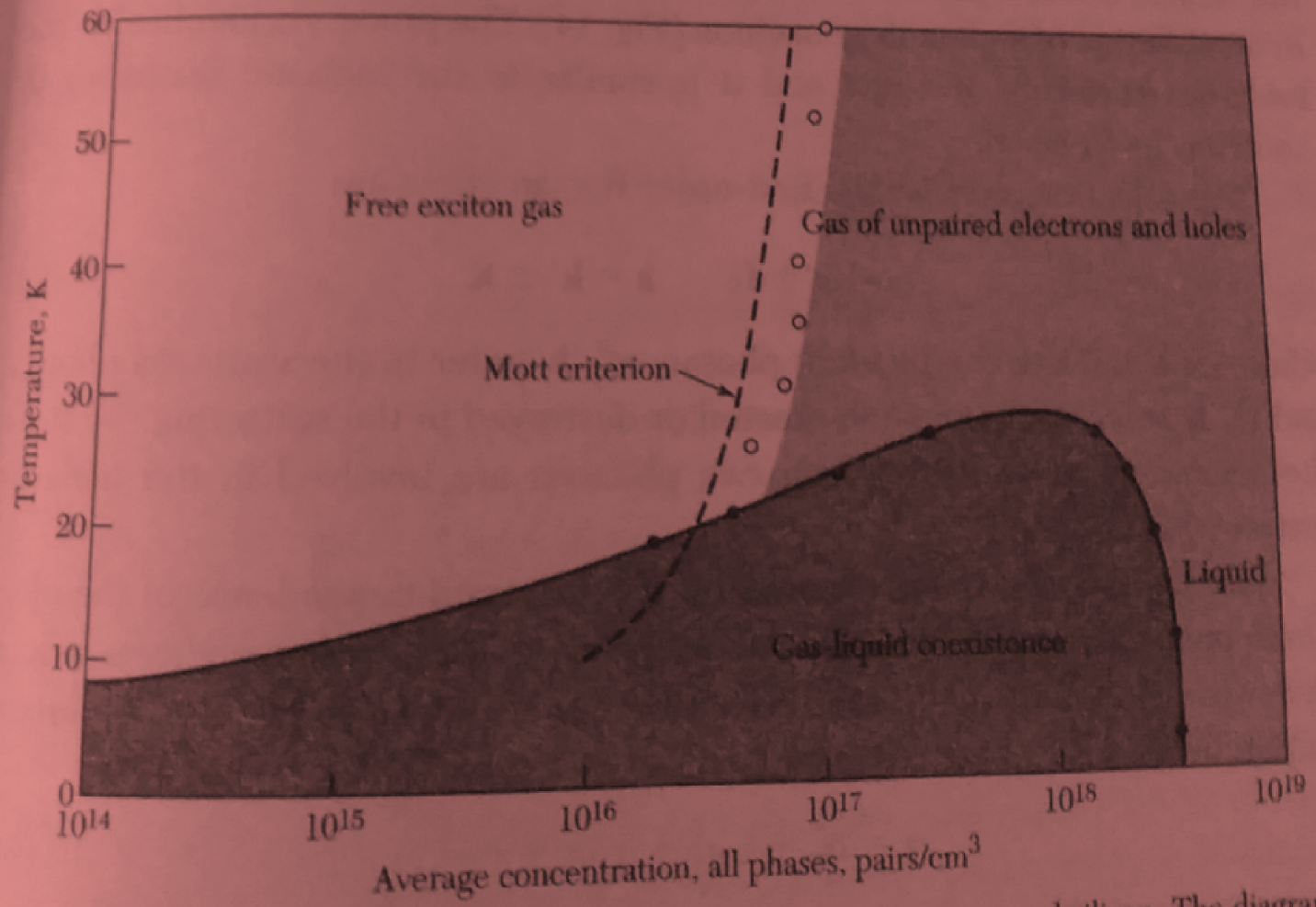


Figure 13 Phase diagram for photoexcited electrons and holes in unstressed silicon. The diagram shows, for example, that with an average concentration near 10^{17} cm^{-3} at 15 K, a free-exciton gas with saturated-gas concentration of 10^{16} cm^{-3} coexists with a (variable) volume of liquid droplets, each with a density of $3 \times 10^{18} \text{ cm}^{-3}$. The liquid critical temperature is about 23 K. Theoretical and experimental values for the metal-insulator transition for excitons are also shown. (From J. P. Wolfe.)

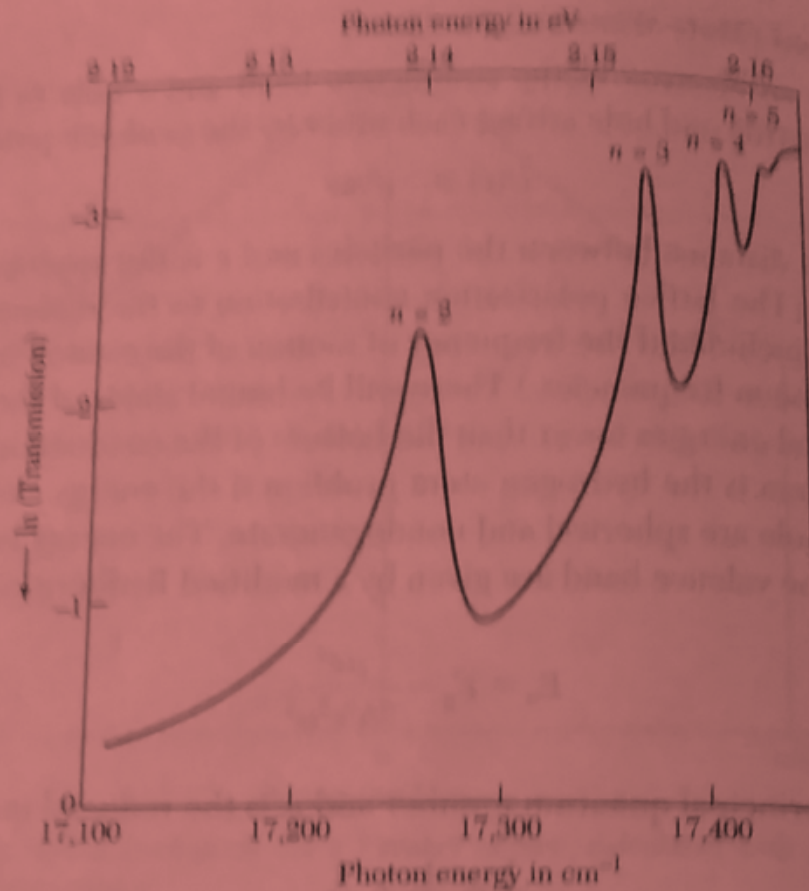


Figure 10 Logarithm of the optical transmission versus photon energy in cuprous oxide at 77 K showing a series of exciton lines. Note that on the vertical axis the logarithm is plotted decreasing upward; thus a peak corresponds to absorption. The band gap E_g is 2.17 eV. (After P. W. Baumeister)

- Within the drop the excitons dissolve into a degenerate fermi gas of electrons and holes , with metallic properties this state was predicted by L.V.Keldysh.
- At low pressure , the exciton gas is insulating .
- At high pressure , the exciton gas breaks up into a conducting plasma of unpaired electron and hole.
- The transition from exciton to the plasma is an example of mott transition.